# Trajectory and orbit of the unique carbonaceous meteorite Flensburg 

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#### Abstract

The C1-ungrouped carbonaceous chondrite Flensburg fell in Germany on September 12,2019 , in the daytime. We determined the atmospheric trajectory, velocity, and heliocentric orbit using one dedicated AllSky6 meteor camera and three casual video records of the bolide. It was found that the meteorite originated in the vicinity of the $5: 2$ resonance with Jupiter at heliocentric distance of 2.82 AU . When combined with the bolide energy reported by the United States government sensors (USGS), the preatmospheric diameter of the meteoroid was estimated to be $2-3 \mathrm{~m}$ and the mass to be $10,000-20,000 \mathrm{~kg}$. The meteoroid fragmented heavily in the atmosphere at heights of $46-37 \mathrm{~km}$, under dynamic pressures of $0.7-2 \mathrm{MPa}$. The recovery of just one meteorite suggests that only a very small part of the original mass reached the ground. The bolide velocity vector was compared with that reported by the USGS. There is good agreement in the radiant but the velocity value has been underestimated by the USGS by almost $1 \mathrm{~km} \mathrm{~s}^{-1}$.


## INTRODUCTION

The Flensburg meteorite fell in northern Germany, close to the Danish border, on September 12, 2019, 12:50 UT (14:50 local daylight saving time). The fall was accompanied by a very bright bolide (a superbolide) easily visible in broad daylight. The International Meteor Organization received 584 reports of visual sightings of the bolide through its online form (https://fireballs.imo.ne $t /$ ). Most of the reports were from the Netherlands, where the weather was best at the time of the event. The most distant sighting, more than 600 km from the bolide, was reported from England.

The bolide was also detected by the spaceborne United States government sensors (USGS, https://cneos. jpl.nasa.gov/fireballs/, see also Brown et al. 2016). The sensors provided the exact time of the bolide (12:49:48 UT), approximate geographic location $\left(54.5^{\circ} \mathrm{N}, 9.5^{\circ} \mathrm{E}\right)$ and height ( 42 km ) of the point of maximum brightness, bolide velocity ( $18.5 \mathrm{~km} \mathrm{~s}^{-1}$ ), and total radiated energy
$\left(1.69 \times 10^{11} \mathrm{~J}\right)$. The total impact energy was deduced to be 0.48 kt TNT , i.e., $2.0 \times 10^{12} \mathrm{~J}$. USGS report only about a dozen bolides with similar or larger energies per year globally. The event was therefore quite significant on a regional scale. The provided velocity components indicated that the bolide moved nearly from the south to the north (heading azimuth $9^{\circ}$ ) on a trajectory with low slope ( $24^{\circ}$ to the horizontal). Note, however, that while the USGS reported energies can be considered reliable, the velocities and trajectories were found to be in severe error for some bolides with independent data available (Borovička et al. 2017; Jenniskens et al. 2018; Devillepoix et al. 2019; Unsalan et al. 2019).

A single meteorite fragment was found by chance the next day, September 13, by Mr. Erik Due-Hansen on the lawn of his yard in Flensburg. The meteorite has been described and analyzed by Bischoff et al. (2021). The total mass was 24.5 g with the bulk density of only $(1984 \pm 16) \mathrm{kg} \mathrm{m}^{-3}$. The meteorite has been classified as C 1 -ungrouped carbonaceous chondrite. It therefore
represents a unique material not found in other known meteorites. No further meteorite fragment from the Flensburg fall was found as of the time of writing (September 2020).

The purpose of this paper is to set the Flensburg meteorite into a geological context within the solar system, i.e., to determine its pre-fall heliocentric orbit and possible source region. At the same time, the atmospheric trajectory of the bolide will be studied with the aim to obtain information about the physical properties of the original meteoroid. Since we cannot rely on the USGS data, we will use four ground-based video records of the bolide. One video was obtained by us and three others were found on the Internet and carefully calibrated. Similar work has been done for three other instrumentally observed carbonaceous meteorite falls: C2-ungrouped Tagish Lake (Brown et al. 2000; Hildebrand et al. 2006) and two CM2 meteorites, Sutter's Mill (Jenniskens et al. 2012) and Maribo (Borovička et al. 2019).

## THE VIDEO RECORDS

At the time of the bolide, three AllSky6 meteor observing camera systems were in operation in Germany. The AllSky6 system (Hankey et al. 2020) consists of six video cameras at each site, which run 24 h a day. As is illustrated on an image from the meteorological satellite Meteosat in Fig. 1, the early afternoon of September 12 was partly or mostly cloudy over the majority of Germany and the surrounding countries. Only the camera in Herford, operated by one of us (JS), recorded part of the bolide, about 1 s long, in a gap between clouds (Fig. 2). Since the camera is stationary, it was easy to calibrate the video astrometrically. Stars were measured on four saved video sequences (showing other meteors) from the night of September 10/11 and three sequences from September 13/14. All data fit well together, confirming that the camera did not move in between. In total, 491 star positions were used. Nine plate constants of the axially symmetric projection described in Borovička (2014) were determined. The plate coordinates $x, y$ could then be transformed into azimuths and elevations $A, h$.

Several casual video records of the bolide were published on the Internet. From these, we selected three videos which were suitable for on-site stellar calibration and showed the bolide from different angles. All three videos were recorded by dashboard cameras in moving cars. The car motion complicated the calibration; nevertheless, we were able to extract the bolide coordinates well, especially from two videos. We also tried to calibrate another video taken by a stationary security camera in Marknesse, the Netherlands, but the
low quality of the camera and lack of sufficient terrestrial calibration objects prevented us from obtaining useful data from that record. The camera scale was different for $x$ - and $y$-axes and we were unable to obtain the correct ratio.

A very useful video was taken on the Energieweg in Nijmegen, the Netherlands. Some images from the video are shown in Fig. 3. The bolide appeared as a small spot near the center of the frame. It brightened slowly before flaring up suddenly. After the flare, the main fragment continued down, leaving a bright trail behind it, visible on several subsequent frames. The whole bolide, lasting for 4.5 s , passed on a clear part of the sky, although a few frames, where the bolide crossed high-voltage aerial power cables, were not measurable. The video frame shown in Fig. 3 was calibrated astrometrically using the method described in Borovička (2014). The nighttime starry images were taken from nearly the same spot by one of us (FB) on December 18, 2019. The left third of the image was to be covered because of too-bright street lights. Still, a sufficient number of stars (38) could be measured. At the same time, 23 terrestrial objects, well visible on both the original video and the calibration image, were measured. Their distances from the calibration camera were measured on Google Earth and ranged from 30 m for nearby objects to almost 2 km for a distant bridge construction. The difference between the position of the calibration camera and the actual position of the original video camera was found by minimizing the residua of the calibration. The difference was about 1.3 m , nearly in the direction of car motion. This terrestrial coordinate correction was taken into account in the final astrometric calibration (see Borovička [2014] for details of the calibration procedure).

Since the car did not move straight all the time, the aiming point of the camera changed. The changes were determined by measuring the distant bridge on each video frame. The plate constants $A_{0}, h_{0}$ (azimuth and elevation of the aiming point) were adjusted accordingly for each frame. Other plate constants were kept fixed. The distance covered by the car during the bolide flight was ignored since it was small in comparison with the scatter of the measurements at the bolide distance (in fact only the projection of the car motion perpendicular to the line of sight may play a role and that was quite small because the car moved toward the bolide).

Another useful video was taken at a road near the northern tip of the German island Rügen. Some frames of this video are presented in Fig. 4. Also here the bolide was located near the center of the field of view, which is convenient because image distortion is negligible here. At the beginning, the bolide emerged from behind the clouds and the bolide end was also


Fig. 1. Image of Germany and surrounding areas taken by the meteorological geostationary satellite Meteosat on September 12, 2019, 12:45 UT (courtesy EUMETSAT and Czech Hydrometeorological Institute). Ground projection of the bolide trajectory (red line) and locations of the four videos used for the trajectory determination are indicated ( $1=$ Herford, $2=$ Nijmegen, 3 = Rügen, 4 = Almere). (Color figure can be viewed at wileyonlinelibrary.com.)
hidden by clouds. Nevertheless, a significant part of the bolide, about 3 s long, was observed, though also that part was interrupted three times by clouds. The calibration image was taken by one of us (GB) on March 11, 2020. In total, 111 stars and 24 terrestrial objects in distances from 30 m to 2.7 km were measured. No correction of camera position was necessary. On the other hand, since the tilt of the video camera changed during the car motion (as the driver was passing a cyclist), three plate constants, $A_{0}, h_{0}$, and $\varphi$, were to be adjusted for each frame using two points on the distant horizon.

Finally, a video taken in Almere, the Netherlands, was calibrated (Fig. 5). The bolide started to be visible
on a clear part of the sky but disappeared behind clouds 1 s later. After another 2 s , the bolide reappeared shortly (on three frames) in a small gap between clouds. Since the car turned on a roundabout in between, both parts of the video were calibrated independently using separate calibration images taken by one of us (FB) on July 11, 2020. There were 76/78 stars and $21 / 18$ terrestrial objects in distances from 10 m to 0.5 km . Corrections of camera positions were applied. Because of lower quality of this video and less convenient terrestrial calibration objects, the calibration of this video was less robust than was the case of the previous two casual videos. A range of possible solutions, varying in the values of camera position


Fig. 2. Co-added video frames from the Herford video, converted to grayscale, showing the bolide path in a gap between clouds. The color inset shows an enlargement from a single frame. Author of the video: Jörg Strunk. (Color figure can be viewed at wile yonlinelibrary.com.)
correction and by including/excluding some of the calibration objects, was obtained. The selected solution is fully consistent with the bolide trajectory and velocity determined from other three videos. Other formally allowable solutions resulted in deviations of lines of sight from the bolide trajectory up to about $\pm 1 \mathrm{~km}$. The bolide trajectory presented in the next section was therefore primarily based on the videos from Herford, Nijmegen, and Rügen. The Almere video demonstrated that there is no conflict with that solution.

The parameters of all four videos including geographical coordinates are given in Table 1. The video locations relative to the bolide path are shown in Fig. 1.

## BOLIDE TRAJECTORY AND VELOCITY

The bolide trajectory was computed by the least squares method of Borovička (1990). Video calibrations provided the source data: bolide azimuths and elevations for each video frame. The source data can be found in File S1 in supporting information. In Herford, the bolide was visible at elevations $17^{\circ}-14^{\circ}$ above the horizon. In all other videos, the first measurements were at elevations between $11^{\circ}$ and $12^{\circ}$. The end of the bolide was visible only from Nijmegen, $3.5^{\circ}$ above the horizon.

The trajectory was first assumed to be straight. Corrections for curvature due to gravity were applied at the end, when the linear trajectory and velocity were known. The least squares method finds the linear trajectory by minimizing the sum of squares of the distances (in space) between the trajectory and the lines of sight. Since the videos and their calibrations were of varying quality, weights were set differently for each video. Herford got the weight 100, Nijmegen and Rügen 1, and Almere 0.1. Nevertheless, the results did not differ much for equal weights: by less than $0.1^{\circ}$ for the radiant and by about 100 m for the position of the end point. Good consistency of all measurements is demonstrated in Fig. 6, where the deviations of lines of sight from the trajectory are plotted. There is no significant systematic trend of any video. The points are randomly mixed. For Herford, most points lie within 50 m of the trajectory. The scatter of the other videos is about $\pm 300 \mathrm{~m}$. This larger scatter is mostly due to the fact that the cameras were located in moving cars.

The observed trajectory was 90 km long. The bolide was first detected at a height of 71.8 km from Almere. From Nijmegen, the first detection was at a height of 70.5 km . The Herford camera covered a stretch 18 km long between heights 67.0 and 59.6 km . The last


Fig. 3. Single frame from the Nijmegen video. The bolide is nearly in the middle, better visible in the inset at upper left, where part of the image was enlarged and contrast was enhanced. The white arrow above the center of the frame indicates where the bolide started to be visible 2.5 s earlier. The other two insets show the bolide in the maximum light and toward the end. Author of the video: Gerard Kemna. The video was downloaded from the site of the Dutch news media NU.nl (https://www.nu.nl/binne nland/5991876/heldere-vallende-ster-op-klaarlichte-dag-waargenomen-boven-nederland.html). (Color figure can be viewed at wile yonlinelibrary.com.)
observation was from Nijmegen at a height of 35.3 km . The average slope of the trajectory to the horizontal was $24.4^{\circ}$ (the slope changes along the trajectory due to the Earth's curvature).

After the trajectory was computed, individual measurements were projected onto it. The dependency of the length along the trajectory (i.e., the distance travelled from the bolide beginning) as a function of time was obtained this way. The relative time was computed for each video from the known frame rate and considering that the cameras had a rolling shutter (i.e., the frames were not read out at once but line by line). These relative times are provided in File S1. The time offsets between cameras were then determined using the bolide data. The following time offsets relative to the Herford camera were found: Nijmegen -0.4549 s , Rügen 0.4121 s , Almere -0.5556 s . Using the time of bolide maximum reported by the USGS, the time zero at Herford was found to correspond to 12:49:45.2 UT.

From the Herford data alone, the average velocity of the bolide in the covered stretch was $19.422 \pm 0.020 \mathrm{~km} \mathrm{~s}^{-1}$. The data from the other videos
are fully consistent with this value but show a larger scatter. In Fig. 7 the difference between the observed length and the expected length at the given time is plotted. Almere and Nijmegen show the largest scatter; Rügen shows smaller deviations. The scatter grows near the end of the trajectory, where the bolide was fragmented and more difficult to measure. Again, there is no systematic trend among the videos. Note that the velocity is not expected to remain constant along the whole trajectory. Atmospheric drag causes deceleration even for bolides of this size. The lag expected from bolide modeling assuming an initial mass of $22,000 \mathrm{~kg}$ and severe fragmentations at lengths 63 and 71 km (heights 45.5 and 42.5 km ) derived from the light curve (see the Mass and Strewn Field section) is indicated by the dotted line in Fig. 7.

Since the flight was rather horizontal, the Earth's gravity must have caused some bending of the trajectory. It can be calculated that after the 4.55 s covered by the observations, gravity will displace the meteoroid by about 90 m in lateral direction. This is less than the scatter of the data (see Fig. 6), so the trajectory curvature could not be directly measured.


Fig. 4. Single frame from the Rügen video. The bolide is nearly in the middle, better visible in the middle inset, where part of the image was enlarged. The black arrow indicates where the bolide started to be visible 3 s earlier. The left inset shows the bolide at maximum brightness, which occurred just six frames earlier. The right inset shows one of the last frames, where the bolide is visible. Author of the video: Holger Scheele. The video was downloaded from https://www.youtube.com/watch?v= W0ZFqBDdTJk. (Color figure can be viewed at wileyonlinelibrary.com.)

Nevertheless, the expected change in the radiant was taken into account when computing the heliocentric orbit.

Trajectory, radiant, and velocity data are provided in Table 2. The beginning and end points are the points where the bolide started and ceased to be visible on the videos, respectively. Should the bolide appear at night, it would certainly have started to be visible at higher height and could be probably followed to somewhat lower height. The errors of the heights quoted in Table 2 express the uncertainty in the position and orientation of the trajectory, not the uncertainty of visibility of start/end point. The point of maximum brightness was not well defined since the bolide exhibited similar brightness at the heights of $43-40 \mathrm{~km}$. The entry speed takes into account the expected slight atmospheric deceleration before the bolide appeared on the Herford video.

For comparison, data from the USGS are given in Table 2 as well. We note that the USGS position is found about 20 km to the east. The radiant agrees within one degree. Note that radiant azimuth and zenith distance change along the trajectory due to the Earth's curvature. Right ascension and declination change only
slightly due to the curvature of the trajectory. The biggest discrepancy to the USGS data is the bolide velocity, which was reported to be almost $1 \mathrm{~km} \mathrm{~s}^{-1}$ lower than precisely deduced in this work. The difference is only partly explained by atmospheric deceleration between the beginning and the maximum brightness point. The measured average velocity between the heights of 45 and 41 km (from the Nijmegen video) was $19.3 \pm 0.5 \mathrm{~km} \mathrm{~s}^{-1}$ and the modeled velocity at 42 km was $19.1 \mathrm{~km} \mathrm{~s}^{-1}$, i.e., still $0.6 \mathrm{~km} \mathrm{~s}^{-1}$ larger than the USGS value.

Figure 8 shows the ground projection of the trajectory on the map. The bolide appeared above the Elbe river estuary and then continued nearly northward, above land.

The geocentric radiant and heliocentric orbit were computed from the apparent radiant and entry velocity by the analytical method of Ceplecha (1987). The results are given in Table 3. The orbit is compared graphically with the orbits of three other carbonaceous chondrites with a known orbit in Fig. 9. The orbit of Flensburg, whose fall was captured with one dedicated meteor camera, is probably the most precise among them while the orbit of the Tagish Lake, which was not directly


Fig. 5. Single frame from the Almere video showing the bolide (in red circle) shortly reappearing between clouds. The white arrow indicates where the bolide started to be visible 3 s earlier. Author of the video: Leon Pepping. The video was downloaded from the site of the Dutch news media omroepflevoland.nl (also available at https://www.youtube.com/watch?v=Qp2_ BZju7EU). (Color figure can be viewed at wileyonlinelibrary.com.)

Table 1. Video records used for trajectory determination.

| No. | Source | Site | Longitude ${ }^{\circ} \mathrm{E}$ | Latitude ${ }^{\circ} \mathrm{N}$ | Altitude $(\mathrm{m})$ | Resolution pixels | FOV $\left({ }^{\circ}\right)$ | FPS | NOF |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 1 | Meteor camera | Herford | 8.70178 | 52.12364 | 121 | $1920 \times 1080$ | $87 \times 47$ | 25 | 23 |
| 2 | Moving car | Nijmegen | 5.82536 | 51.84390 | 8 | $1920 \times 1080$ | $68 \times 40$ | 25 | 106 |
| 3 | Moving car | Rügen | 13.33233 | 54.65956 | 17 | $1920 \times 1080$ | $95 \times 58$ | 30 | 56 |
| 4 | Moving car | Almere | 5.26111 | 52.40027 | 0 | $960 \times 540$ | $46 \times 25$ | 25 | 24 |

The last three columns contain the estimated field of view (FOV), the number of frames per second (FPS), and the number of frames with visible bolide (NOF).
imaged from the ground, is the least reliable. All orbits lie inside the orbit of Jupiter. The inclinations are small in all cases. The inclination of Flensburg $\left(6.8^{\circ}\right)$ is the largest of all four. Also the aphelion distance of Flensburg ( 4.8 AU ) is the largest one but is comparable with Maribo (4.4 AU) and Sutter's Mill (4.7 AU). These two CM2 chondrites have lower perihelia, higher eccentricities, and lower semimajor axes than Flensburg. The most significant aspect of the Flensburg orbit is that it was in $5: 2$ resonance with Jupiter. This powerful resonance causes one of the Kirkwood gaps in the main asteroid belt, where all orbits become chaotic (e.g., Moons 1996). Although the Tisserand parameter of 2.89 with respect to Jupiter may suggest a cometary orbit, the fact that it was in the $5: 2$ resonance allows us to
classify the Flensburg orbit as asteroidal (Tancredi 2014).

## LIGHT CURVE AND FRAGMENTATION

Since the fall occurred during the daytime, a detailed radiometric light curve, such as the case for the Maribo fall (Borovička et al. 2019), is not available. To get at least an approximate light curve, we used the Nijmegen video. Assuming that the camera has a gamma factor of 0.45 , the bolide signal on each video frame was measured. An ad hoc correction for signal saturation was applied for the period of large brightness. Finally, the light curve was scaled to get the total radiated energy of $1.69 \times 10^{11} \mathrm{~J}$ as reported by


Fig. 6. Deviations of individual lines of sight from four video sites from the bolide trajectory. The deviation is positive if the line of sight passes above the trajectory or, in the case of Herford, which lay almost in the fall plane, east of it. Note the different scales on the $x$ - and $y$-axes. (Color figure can be viewed at wileyonlinelibrary.com.)


Fig. 7. Deviations of the measured lengths along the trajectory from the length expected for the given time and a constant bolide velocity of $19.42 \mathrm{~km} \mathrm{~s}^{-1}$. The dotted line shows the expected deceleration. (Color figure can be viewed at wileyonlinelibrary.com.)
the USGS. The resulting light curve is presented in Fig. 10. It must be kept in mind that the way to obtain it was not rigorous. Because the values of absolute magnitude and the exact shape of the light curve could not be determined independently, their reliability is limited. Nevertheless, the main characteristics are obvious. The brightness was increasing steadily from the beginning down to the height of 45 km , where the bolide exploded, i.e., the brightness increased abruptly. The increase was followed by a broad maximum at heights $45-40 \mathrm{~km}$. Thereafter, the brightness decreased, possibly with some minor flares on the descending branch.

Table 2. Trajectory, apparent radiant and velocity, including comparison with the USGS data.

|  | This work | USGS |
| :--- | :--- | :--- |
| Beginning point |  |  |
| Longitude $\left({ }^{\circ} \mathrm{E}\right)$ | $9.0322 \pm 0.0006$ |  |
| Latitude $\left({ }^{\circ} \mathrm{N}\right)$ | $53.887 \pm 0.002$ |  |
| Height $(\mathrm{km})$ | $71.84 \pm 0.08$ |  |
| Maximum brightness |  | 9.5 |
| Longitude $\left({ }^{\circ} \mathrm{E}\right)$ | 54.492 | 54.5 |
| Latitude $\left({ }^{\circ} \mathrm{N}\right)$ | 40.7 | 42 |
| Height $(\mathrm{km})$ |  |  |
| End point | $9.2060 \pm 0.0016$ |  |
| Longitude $\left({ }^{\circ} \mathrm{E}\right)$ | $54.598 \pm 0.002$ |  |
| Latitude $\left({ }^{\circ} \mathrm{N}\right)$ | $35.3 \pm 0.08$ | 188.8 |
| Height $(\mathrm{km})$ |  |  |
| Apparent radiant* | $188.1 \pm 0.1$ | 66.4 |
| Azimuth $\left({ }^{\circ}\right)$ | $65.3 \pm 0.1$ | 184.7 |
| Zenith distance $\left({ }^{\circ}\right)$ | $185.00 \pm 0.10$ | -11.4 |
| Right ascension $\left({ }^{\circ}\right)$ | $-11.01 \pm 0.12$ | 18.5 |
| Declination $\left({ }^{\circ}\right)$ | $19.43 \pm 0.05$ |  |
| Entry velocity* $\left(\mathrm{km} \mathrm{s}{ }^{-1}\right)$ |  |  |

*Apparent radiant is given at the beginning of the trajectory. Entry velocity is given at the top of the atmosphere. USGS data are valid for the maximum brightness point. Azimuth is counted from the north clockwise. Equatorial coordinates are for the equinox J2000.0.

The light curve is very different from that of Maribo, which exhibited a steep increase at the beginning and many relatively short flares in the second half of the trajectory (Borovička et al. 2019). Maribo fragmented many times, which was interpreted as an evidence for the inhomogeneous nature of the meteoroid. The meteoroid contained parts with various strengths, which disintegrated at various dynamic pressures. The initial disruption of Maribo occurred at 0.017 MPa and other fragmentations occurred at $0.25-$ 4.3 MPa .

The main disruption of Flensburg started at 0.7 MPa and continued until about 2 MPa . The interval of dynamic pressures causing fragmentation was therefore much smaller than that of Maribo. Judging from the light curve, a small fragmentation may have occurred also at the beginning of the bolide (where the brightness increased above the detection limit), i.e., at a height of $\sim 71 \mathrm{~km}$ and a dynamic pressure of $\sim 0.025 \mathrm{MPa}$, but only a small amount of mass ( $\sim 1 \%$ ) may be lost there. We conclude that the Flensburg meteoroid was therefore more homogeneous than Maribo. In this respect it more resembles the Romanian superbolide of January 7, 2015 (Borovička et al. 2017), where, however, no meteorites were recovered and the whole body was probably completely pulverized in the atmosphere.

Additional information about meteoroid fragmentation can be obtained from bolide images.


Fig. 8. Ground projection of the bolide trajectory and the position of the meteorite (marked M). The source of the background image is Google Earth. (Color figure can be viewed at wileyonlinelibrary.com.)

Table 3. Geocentric radiant and heliocentric orbit (J2000.0).

| Geocentric right ascension of radiant | $\alpha_{\mathrm{g}}$ | $183.46^{\circ} \pm 0.11^{\circ}$ |
| :---: | :---: | :---: |
| Geocentric declination of radiant | ¢g | $-18.18^{\circ} \pm 0.14^{\circ}$ |
| Geocentric entry speed ( $\mathrm{km} \mathrm{s}^{-1}$ ) | $v_{\mathrm{g}}$ | $15.97 \pm 0.06$ |
| Semimajor axis (AU) | $a$ | $2.82 \pm 0.03$ |
| Eccentricity | $e$ | $0.701 \pm 0.003$ |
| Perihelion distance (AU) | $q$ | $0.843 \pm 0.001$ |
| Argument of perihelion | $\omega$ | $307.25^{\circ} \pm 0.16^{\circ}$ |
| Longitude of ascending node | $\Omega$ | $349.207^{\circ} \pm 0.001^{\circ}$ |
| Inclination | $i$ | $6.82^{\circ} \pm 0.06^{\circ}$ |
| Aphelion distance (AU) | $Q$ | $4.80 \pm 0.06$ |
| Perihelion date |  | $\begin{aligned} & \text { 2019-08- } \\ & 07.6 \pm 0.1 \text { day } \end{aligned}$ |
| Tisserand parameter | $T$ | $2.89 \pm 0.02$ |
| Orbital period | P | $4.74 \pm 0.08 \mathrm{yr}$ |
|  | $\frac{P_{\text {jup }}}{P}$ | $2.502 \pm 0.007$ |

Figure 11 shows selected frames from a casual video taken during a boat trip to the Wangerooge Island, i.e., relatively close to the bolide $(\sim 120 \mathrm{~km})$. The video could not be used for trajectory determination because it was taken from the sea on an unstable boat with no fixed reference objects. Moreover, the exact position of the boat is unknown. The video shows the bolide before the main flare, at heights of about $54-45 \mathrm{~km}$, including the onset of the flare. During the maximum phase, the bolide was hidden behind a boat passenger and emerged again when it was at about 39 km . The heights were determined from the relative time and the time-height relation known from other videos. The video time was related to the other videos using the onset of the main flare.

As can be seen in Fig. 11, the bolide was an elongated object before the main flare. It means that some smaller fragments were trailing the main body. The initial fragmentation may have been therefore more severe than it seems from the light curve. The main flare started with brightening of the bolide head, i.e.,


Fig. 9. Known orbits of carbonaceous chondrites in projection onto ecliptic plane: Tagish Lake (C2 ungrouped, fell in 2000, orbital data from Hildebrand et al. 2006), Maribo (CM2, 2009; Borovička et al. 2019), Sutter's Mill (CM2, 2012; Jenniskens et al. 2012), and Flensburg (C1 ungrouped, 2019; this work). (Color figure can be viewed at wileyonlinelibrary.com.)
fragmentation of the main body. When the bolide again emerged on the video after the maximum, it had a bright head and a long and bright wake. After a short time, the head faded and separated from the wake. During the final phase, no forward motion was apparent any more. Instead, two separated sections of the stationary trail faded gradually. This behavior suggests that no large fragments were formed and the meteoroid disintegrated into small (subkilogram) fragments and dust. However, because of the daylight conditions and thus limited sensitivity of the camera, the presence of an individual large fragment $(\sim 10 \mathrm{~kg}$, see also the next section) at the end cannot be strictly excluded.

## MASS AND STREWN FIELD

Using the total impact energy of $2 \times 10^{12} \mathrm{~J}$ reported by the USGS and the known initial velocity of $19.43 \mathrm{~km} \mathrm{~s}^{-1}$, the meteoroid initial mass can be calculated to $10,000 \mathrm{~kg}$. If the light curve in Fig. 10 is fitted by the semiempirical model in a similar way as Maribo (Borovička et al. 2019), the value of $22,000 \mathrm{~kg}$ is obtained from the (generally uncertain) brightness before the main flare. Taking these two values as the possible mass range and assuming bulk density between 1600 and $2000 \mathrm{~kg} \mathrm{~m}^{-3}$ (the meteorite density was $1984 \mathrm{~kg} \mathrm{~m}^{-3}$ ), the resulting initial equivalent diameter of the meteoroid is between 2 and 3 m .


Fig. 10. Approximate light curve of the Flensburg bolide based on the Nijmegen video.

The available data are not sufficient to rigorously determine the total fallen mass and/or the number of meteorites. Since the fragmentation was enormous and the bolide terminated at a relatively large height of 35 km , the recovered meteorite of 24.5 g may belong to the largest meteorites which reached the ground. Then, taking into account the large initial mass and assuming a power-law fragment mass distribution with mass index 2 and upper mass limits of $10-100 \mathrm{~g}$ in individual fragmentations, hundreds of meteorites of about 10 g and tens of thousands of about 1 g may exist.

Dark flight computation was performed to find the expected spread of the meteorites on the ground. Meteorite fragments were assumed to be formed at two fragmentation points at the heights of 45.5 and 42.5 km , forming the wide bolide maximum, and at a final fragmentation of the remaining piece at the height of 37 km , forming a smaller flare on the light curve (see Fig. 10). The ablation of the formed fragments was computed until their velocities dropped to $2.5 \mathrm{~km} \mathrm{~s}^{-1}$. The ablation coefficient of $0.005 \mathrm{~s}^{2} \mathrm{~km}^{-2}$ was used; the resulting strewn field is, nevertheless, not sensitive to the exact value of the ablation coefficient. Meteoroid density was assumed to be $2000 \mathrm{~kg} \mathrm{~m}^{-3}$ and the value of $Г А$ (product of drag coefficient and shape coefficient) was assumed to be 0.8 . Other values of the density and ГА would produce an identical strewn field but the meteorite masses at a given location would be different. The high altitude winds and other atmospheric parameters for dark flight computation were taken from two numerical weather prediction models, European Centre for Medium-Range Weather Forecasts (ECMWF) and Aire Limitée Adaptation dynamique Développement InterNational (ALADIN). Both models


Fig. 11. Bolide images from a video taken at a boat to Wangerooge Island. The six main images show the final phase of the bolide in 0.1 s intervals. The bolide head is at a height of about 38.5 km in the first image (at the top). The dark object is a shoulder of a boat passenger. The two insets show earlier phases of the bolide using the same scale, at heights of about 51 and 47 km , respectively. Author of the video: The Oldenburg Red Cross. The video was downloaded from the site of The Weather Network (www.theweathernetwork.com; also available at https://www.youtube.com/watch? $\mathrm{v}=30 \mathrm{XQhLKOa}$ 9 k ). (Color figure can be viewed at wileyonlinelibrary.com.)
gave similar results. ECMWF resulted in meteorite positions $300-400 \mathrm{~m}$ more to the north than ALADIN. This value is small in comparison with the expected meteorite spread and, moreover, the difference lies in the same direction as the general orientation of the strewn field, so there were no practical consequences for meteorite searches.

The expected strewn field for meteorite masses 5100 g is shown in Fig. 12. The strewn field is oriented from the south to the north. Larger meteorites lie more to the north, a consequence of the direction of bolide flight. The direction of the strongest winds was from WNW, so the whole strewn field was shifted to the east from the prolonged bolide trajectory. Meteorites originating at the lower height of 37 km are shifted to the north (more exactly, in the azimuth of bolide flight) in respect of the meteorites of the same mass originating higher, simply because they started the dark flight later. As a consequence, the strewn field of meteorites originating at 37 km lies more to the east as of those originating higher, as shown in Fig. 12. When moving from west to east, one should first encounter larger meteorites originating at higher heights (if they exist) and then smaller meteorites from lower heights.

The line for meteorites originating at lower heights passes very close to the position of the recovered meteorite. Although meteorites of masses of about 45 g were predicted there and the actual meteorite had only 24.5 g , that discrepancy can be easily removed by adjusting ГА to 0.65 instead of 0.8 for this meteorite, which is quite plausible. For a perfect sphere, the value of $\Gamma$ A would be $\sim 0.6(A=1.21, \Gamma \sim 0.5)$. We can therefore conclude with a high degree of reliability that the meteorite originated from a part of the meteoroid, which survived the main disruption at heights above 40 km and fragmented at lower heights around 37 km at a dynamic pressure about 2 MPa .

Rather extensive searches for more meteorite fragments remained unsuccessful. First sporadic searches around the position of the first meteorite started 10 days after the find. More systematic searches were performed in the second half of November 2019


Fig. 12. Map of the expected meteorite fall area near Flensburg. North is up. The position of the only recovered meteorite is marked M. The modeled meteorite position is depicted by two orange lines. Meteorites originating in the main fragmentation event at the heights of 43-46 km should be spread along the left (western) line. Meteorite masses increase from the south to the north. The extent corresponds to meteorite masses from about 5 to 100 g , assuming a meteorite density of $2000 \mathrm{~kg} \mathrm{~m}^{-3}$ and $\Gamma \mathrm{A}=0.8$. Meteorites of masses $5-50 \mathrm{~g}$ originating in the last major fragmentation event at a height of $\sim 37 \mathrm{~km}$ are expected to be spread along the right line. The source of the background image is Google Earth. (Color figure can be viewed at wileyonlinelibrary.com.)
and then in February 2020 after the bolide trajectory was computed from the first two calibrated videos.

## DISCUSSION

Although the available bolide records were not ideal, the obtained trajectory and velocity are reliable and the precision is reasonably good. In addition to clarifying the origin of the meteorite, this result also represents one of the few opportunities to compare the velocity vector reported by the USGS with independent data. Any details how the USGS data are obtained have not been published. As the USGS dataset is being used to infer properties of meter-scaled meteoroids (Brown et al. 2016; Wiegert et al. 2020), it is important to know the reliability of the data. In this particular case, the direction of flight (the radiant) was determined well (within one degree). The velocity was by $0.9 \mathrm{~km} \mathrm{~s}^{-1}$ lower than the actual entry velocity. The height of the maximum brightness point was determined well. The geographic position was shifted by $0.3^{\circ}$ in longitude $(\sim 20 \mathrm{~km})$. Overall, the Flensburg bolide belongs to the best observed bolides by the USGS (cf. Devillepoix et al. 2019). Still, if USGS data were used to compute the heliocentric orbit, the resulting eccentricity ( 0.64 ), and thus also semimajor axis (2.37 AU) and aphelion distance (3.9 AU) would be too small in comparison with the results derived in this work. The perihelion distance and all angular elements would be nearly identical.

The most important fact revealed by this study is that the Flensburg meteoroid was in the 5:2 resonance with Jupiter located at 2.82 AU . Since the resonance is able to quickly raise the orbital eccentricity and thus bring the meteoroid into an Earth crossing orbit, it explains the extremely short cosmic ray exposure age (CRE) of about 7 kyr (Bischoff et al. 2021). The parent asteroid of Flensburg was very probably orbiting close to the inner or outer edge of the resonance, so that Flensburg entered the resonance soon after being ejected by a collisional event.

Granvik and Brown (2018) evaluated the most likely source regions of 25 meteorites with known orbits. From them, only the H chondrites Košice and Ejby (for more details about their falls, see Borovička et al. [2013] and Spurný et al. [2017], respectively) had more than $10 \%$ probability of coming through the $5: 2$ resonance $(\sim 18 \%$ and $25 \%$, respectively). Košice, however, had a higher probability of coming from the 3:1 resonance located closer to the Sun (at 2.5 AU ). For Ejby, the highest probability ( $\sim 55 \%$ ) was obtained for coming from a Jupiter family comet. Since that origin can be effectively excluded for an ordinary chondrite, 5:2 resonance is the most likely source origin of Ejby, although the $3: 1$ resonance with computed probability
of $10-15 \%$ cannot be fully excluded. Aside from Ejby and Košice, the CM2 chondrites Maribo and Sutter's Mill were the only two meteorites with non-negligible probability of cometary origin on the basis of their orbits. But the $3: 1$ and $v 6$ resonances were nearly equally possible in both cases, while the 5:2 resonance was unlikely ( $\sim 5 \%$ ). Jenniskens (2020) attributed CM chondrites to the $3: 1$ resonance. Flensburg is therefore the only meteorite to date clearly coming through the 5:2 resonance.

In contrast to its young CRE, Flensburg contains very old carbonates, not affected since their formation 4564 Myr ago (Bischoff et al. 2021). This fact may be related to the location of its parent asteroid farther from the Sun, where collisions are less frequent and less energetic, than for most other meteorites.

A strange aspect is the recovery of only a single meteorite. With the initial mass of $10,000-20,000 \mathrm{~kg}$, Flensburg was comparable to the carbonaceous chondrites Tagish Lake ( $50,000-90,000 \mathrm{~kg}$; Brown et al. 2002; Hildebrand et al. 2006) and Sutter's Mill (20,000$80,000 \mathrm{~kg}$; Jenniskens et al. 2012), which both yielded much more meteorites. One possible explanation can be that the recovered meteorite represents an exceptionally compact part of the original meteoroid and the number of fallen meteorites of substantial mass ( $>1 \mathrm{~g}$ ) was indeed low. The fact that the vast majority of radiation was released above the height of 39 km suggests that $\sim 3 \%$ of the original mass reached the height of 37 km , where the meteorite originated, in the form of one or more large bodies ( $\gg 10 \mathrm{~kg}$ ). The fragmentations at higher altitudes may have produced only dust or meteorites of negligible mass. Most of the small fragments formed at 37 km may have crumbled repeatedly, even during the dark flight, so only few of them reached the ground. Unfortunately, the available data are not sufficient to confirm this scenario. In terms of maximum dynamic pressure, the 2 MPa reached by Flensburg is somewhat lower than for Maribo ( 5 MPa ; Borovička et al. 2019) but similar to Tagish Lake (Brown et al. 2002; Popova et al. 2011). Flensburg was therefore not exceptional in its fragmentation strength. Perhaps the more important question is what the outcome of fragmentation is. The Romanian superbolide reached up to 3 MPa and produced no meteorites (Borovička et al. 2017).

Another possibility is that Flensburg produced many meteorites but they degraded quickly in wet European climate. The first meteorite was found the next day after the fall but thorough systematic searches were delayed by 2 months. Note that Tagish Lake meteorites were preserved in ice (Hildebrand et al. 2006), all 78 Sutter's Mill meteorites were found in the two months following their fall (Jenniskens et al. 2012),
and other carbonaceous chondrites were typically found either soon or in drier environments.

## CONCLUSIONS

Flensburg is a unique carbonaceous meteorite. It was therefore important to determine its pre-fall heliocentric orbit. The results showed clearly that the meteorite originated in the vicinity of the $5: 2$ resonance with Jupiter at a heliocentric distance of 2.82 AU. It is in agreement with its very short CRE, which demonstrates that no slow drift mechanism such as the Yarkovsky effect was involved in the transport of the meteorite to the Earth. Flensburg seems to originate from a larger heliocentric distance than other meteorites with a known orbit. The recovery of a single meteorite suggests the ratio of fallen to initial mass was quite low even in comparison with other carbonaceous chondrites. The recovery was enabled by the combination of a large initial mass and the fall into an inhabited area. This fortunate coincidence may explain why no other meteorite of this kind is known.

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## SUPPORTING INFORMATION

Additional supporting information may be found in the online version of this article.

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File S1. Flensburg.xlsx contains bolide celestial coordinates (azimuths and elevations) and relative times at individual video frames. Each of the four videos is given in a separate list.

